

Fig. 2 Temperature dependent UV-vis spectra of L-1 in *n*-dodecane (4.2×10^{-5} M) at 25, 40, 70 and 100 °C. Arrows indicate the change that occurs with increasing temperature. The UV-vis spectrum in TCE at 25 °C is also shown (thin solid line).

In haloalkanes, L-1 shows no CD in the UV-vis range (Fig. S2 and S3[†]). This indicates that the molecules do not form the predominantly preferred chiral structures. L-1, however, shows three major strong Cotton effects with a negative first Cotton effect ($\Delta\epsilon = -100 \text{ M}^{-1} \text{ cm}^{-1}$ at 364 nm), a negative second Cotton effect ($\Delta\epsilon = -270 \text{ M}^{-1} \text{ cm}^{-1}$ at 305 nm) and a positive third Cotton effect ($\Delta\epsilon = 89 \text{ M}^{-1} \text{ cm}^{-1}$ at 292 nm) at 25 °C in *n*-hexane, and similarly in *n*-dodecane (first: $\Delta\epsilon = -108 \text{ M}^{-1} \text{ cm}^{-1}$ at 365 nm; second: $\Delta\epsilon = -266 \text{ M}^{-1} \text{ cm}^{-1}$ at 307 nm; and third: $\Delta\epsilon = 85 \text{ M}^{-1} \text{ cm}^{-1}$ at 293 nm). The slight difference in the CD spectra in *n*-alkanes could be due to a minor change in solvophobic interactions. In addition, the CD spectrum of the other enantiomer, D-1, in *n*-dodecane was the mirror image of that of L-1 as shown in Fig. S2.[†] These results indicate that L-1 and D-1 form opposite preferred-handed helical structures in nonpolar *n*-alkanes.

In a *n*-dodecane solution of L-1 (4.20×10^{-4} M), an NH stretching vibration of amide A was observed at 3262 cm^{-1} ; a similar wavenumber was observed for a solid film of L-1 on KBr (3257 cm^{-1}), whereas a higher wavenumber was observed (3433 cm^{-1}) for a TCE solution of L-1. A similar trend was also observed for the C=O stretching vibration of amide I; 1634 cm^{-1} in *n*-dodecane, 1633 cm^{-1} in a solid film, and 1655 cm^{-1} in TCE. The combined results of the UV-vis, CD and IR indicate that L-1 forms a rigid helical columnar assembly in nonpolar solvents, which is stabilized by strong hydrogen bonds among the large disk cores, while the molecules in haloalkanes are mostly in a molecularly dispersed state.

To investigate the stability of the helical structure, temperature and concentration dependent UV-vis and CD spectroscopy were performed from 25–100 °C (4.20×10^{-5} M) and from 4.20×10^{-4} – 4.20×10^{-7} M (at 25 °C) in *n*-dodecane solutions of 1. A gradual decrease in intensity was observed with increasing temperature along with a small blue shift of the

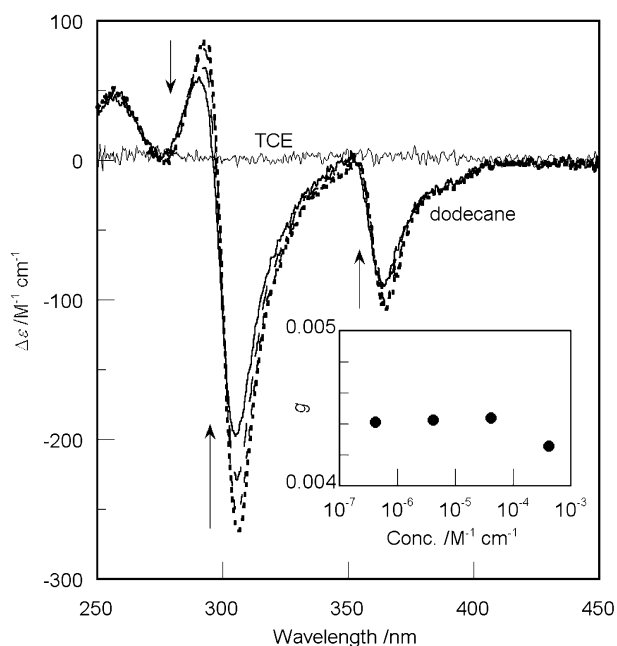


Fig. 3 Temperature dependent CD spectra of L-1 in *n*-dodecane (4.2×10^{-5} M) at 25, 40, 70 and 100 °C. Arrows indicate the change that occurs with increasing temperature. The CD spectrum of L-1 in TCE at 25 °C is also shown (thin solid line). The inset shows the concentration dependence of the chiral anisotropy *g* factor of D-1 in *n*-dodecane at 25 °C.

absorption maximum from 300 to 298 nm (Fig. 2). Correspondingly, there was a gradual decrease in CD intensity with increasing temperature (Fig. 3), which may be attributed to a decrease in positional order of the molecules within the

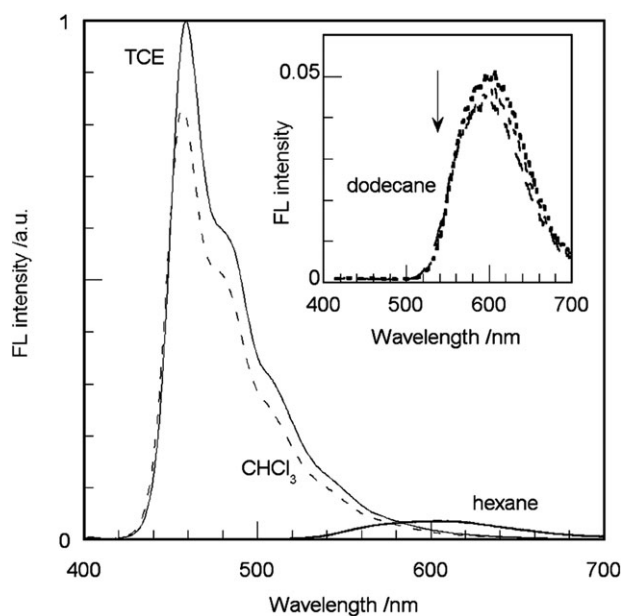


Fig. 4 FL spectra of L-1 in CHCl_3 , TCE and *n*-hexane (4.2×10^{-7} M) at 25 °C. The inset shows an expanded view of the temperature dependent FL spectra in *n*-dodecane at 25, 40 and 70 °C. The arrow indicates the change that occurs with increasing temperature. FL spectra ($\lambda_{\text{ex}} = 360 \text{ nm}$) were normalized using the absorption values at 360 nm.

columns as a result of thermal motion. Importantly, the apparent CD spectrum at 100 °C reveals that the helical structure of **1** is exceptionally thermally stable. Subsequent cooling to 25 °C restored the original CD spectrum. In addition, the chiral anisotropy g factor ($\Delta\epsilon_{307}/\epsilon_{300} = 0.0044$) scarcely changes from low concentrations (4.20×10^{-7} M) to higher concentrations (4.20×10^{-4} M; Fig. 3, inset), suggesting strong association among the molecules.

To further investigate the differences in the intermolecular interactions, fluorescence (FL) spectra (4.20×10^{-7} M) were measured at an excitation of 360 nm (Fig. 4). Emission maxima in CHCl_3 and TCE solutions at 25 °C were observed at 457 nm (fluorescence quantum yield, $\phi = 0.26$) and 459 nm ($\phi = 0.32$), respectively. In contrast to the blue fluorescence observed for these nonhelicogenic solutions, the emission in helicogenic n -hexane and n -dodecane solutions was orange, and their maxima were observed at 604 ($\phi = 0.020$) and 607 nm ($\phi = 0.036$), respectively. For the helicogenic solutions, the broad peaks at around 605 nm are approximately 1/9–1/16 the intensity and are considerably red-shifted with respect to the emission peaks at around 458 nm for the nonhelicogenic solutions. This behavior also indicates that the molecules form columnar structures by strong π - π stacking in the helicogenic solvents.⁹ In addition, close stacking of **1** in n -alkanes is also supported from the fact that the non-aggregated hexakis[4-(hexyloxy-carbonyl)phenylethynyl]benzene, **2**, which does not include amide groups, shows no solvatochromism (Fig. S3†) and its UV-vis and FL spectra in n -hexane were essentially the same as those of **L-1** in nonhelicogenic solutions.¹⁰ Consequently, the FL property is clearly dependent on the aggregation state of **1**, and the emission changes from blue in the molecularly dispersed state to orange in the helical columnar state. Here, the broad emission in the helical columnar state is shifted by approximately 5300 cm^{-1} to lower energy compared to that of monomer emission, considered to be excimer emission, because this energy difference is comparable to those of other excimers generated among the phenylethynyl units.¹¹ With heating, the excimer emission peak hypsochromically shifts by about 4 nm and the intensity gradually decreases; the monomer emission was not observed at all (Fig. 4, inset). This result also supports the strong association of **1** in n -alkane. The decrease in fluorescence intensity is probably due to several factors in combination: (1) a decrease in positional order of the molecules within the columns by thermal motion, (2) breaking of larger columnar stacks into smaller stacks but not monomers,^{12a} and (3) thermally activated nonradiative processes occurring *via* the long chain alkyl group.^{12b}

In summary, we have succeeded in the synthesis of a C_6 -symmetric hydrogen-bonded discotic molecule. Stacking among large central cores of **1** is enforced by six intermolecular hydrogen bonds, resulting in a highly stable supramolecular helical structure, which exhibits excimer emission due to strong intermolecular interactions. In addition, the type of intermolecular association is dependent on the solvent, and the observed fluorescence changes from orange in a helical assembly state to blue in a molecularly dispersed state. Detailed investigations into aspects such as the fluorescence life-time, the self-assembled structure and the chiral amplification,^{4e,f,5b,6a,b,9b} will be performed to further understand features of the supramolecular helicity.

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